

## Locally two-sided approximate solutions in parabolic problems\*

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By means of the fictitious regions method [1] locally two-sided estimations the initial boundary value problem with the second order parabolic equation of the first type are obtained.

In order to find approximate solutions of mathematical physics boundary value problems it is useful to obtain locally two-sided approximate solutions because with their help the approximation order can be increased. By means of the fictitious regions method (FRM) locally two-sided approximate solutions for boundary value problems with the second order elliptic operator are examined in ([2], [3]).

In the present work one type of the auxiliary problem for the second order parabolic equation is determined by means of the FRM. It was proved that two continuations with different sign coefficients of the auxiliary problem are possible. It was shown that the main members of errors in the expansions of the solutions with different signs are also of different signs. And therefore the semi-sum of solutions approximates an exact solution with the order higher than that of each one-sided approximate solution.

Thus, let  $D \subset R^n$  be a finite one-connected domain with the boundary  $\gamma \in C^2$ ,  $T > 0$ , and let  $D_T = \{(x, t) : x \in D, t \in (0, T)\}$  be a cylinder with the side surface  $\gamma_T = \{(x, t) : x \in \gamma, t \in (0, T)\}$ .

Let us consider the initial boundary value problem for the second order parabolic operator of the first kind

$$\begin{aligned} Lu &\equiv \frac{\partial u}{\partial t} - \sum_{i,j=1}^n \frac{\partial}{\partial x_i} \left( a_{i,j}(x, t) \frac{\partial u}{\partial x_j} \right) + a(x, t)u = f(x, t), \quad (x, t) \in D_T, \\ u(x, 0) &= \varphi(x), \quad x \in D, \\ u(x, t) &= 0, \quad (x, t) \in \gamma_T. \end{aligned} \tag{1}$$

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The coefficients  $a_{i,j}(x, t)$  and  $a(x, t)$  satisfy the following conditions:

$$\begin{aligned} a(x, t) &\geq 0, a_{i,j}(x, t) = a_{j,i}(x, t), \\ \nu_0 \sum_{i=1}^n \xi_i^2 &\leq \sum_{i,j=1}^n a_{i,j}(x, t) \xi_i \xi_j \leq \nu_1 \sum_{i=1}^n \xi_i^2, \quad \nu_0 > 0, \end{aligned}$$

for all  $(x, t) \in D_T$  and  $(\xi_1, \dots, \xi_n)$ .

We suppose that

$$\begin{aligned} a_{i,j}(x, t) &\in C(\bar{D}_T), \quad \partial a_{i,j}/\partial x_k(x, t) \in C(\bar{D}_T), \quad a(x, t) \in C(\bar{D}_T), \\ f(x, t) &\in L_2(D_T), \quad \varphi(x) \in \dot{W}_2^1(D). \end{aligned}$$

With such suppositions problem (1) has the only solution which belongs to the space  $W_{2,0}^{2,1}(D_T)$ . In this case the following estimate takes place ([4], p. 388):

$$\|u\|_{W_{2,0}^{2,1}(D_T)} \leq C_1(\|f\|_{L_2(D_T)} + \|\varphi\|_{\dot{W}_2^1(D)}). \quad (2)$$

By means of the fictitious region  $D_1$  we shall supplement  $D$  to some other region  $D_0$  with the boundary  $\Gamma$ , which does not have common points with  $\gamma$ . Let us designate  $D_T^i = \{(x, t) : x \in D_i, t \in (0, T)\}$ ,  $i = 0, 1$ ,  $\Gamma_T = \{(x, t) : x \in \Gamma, t \in (0, T)\}$ .

The auxiliary problem of the FRM is determined in the following way:

$$\begin{aligned} L_\varepsilon u_\varepsilon &= f, \quad (x, t) \in D_T, \quad \frac{\partial u_\varepsilon}{\partial t} - \Delta u_\varepsilon = 0, \quad (x, t) \in D_T^1, \\ u_\varepsilon(x, 0) &= \varphi, \quad x \in D, \quad u_\varepsilon(x, 0) = 0, \quad x \in D_1, \\ u_\varepsilon(x, t) &= 0, \quad (x, t) \in \Gamma_T; \end{aligned} \quad (3)$$

On the surface  $\gamma_T$ , where coefficients are discontinuous, the agreement conditions are assumed

$$u_\varepsilon|_{\gamma_T}^+ = u_\varepsilon|_{\gamma_T}^-, \quad \frac{\partial u_\varepsilon}{\partial N}|_{\gamma_T}^+ = \frac{Q}{\varepsilon} \frac{\partial u_\varepsilon}{\partial n}|_{\gamma_T}^-. \quad (4)$$

Here  $\varepsilon > 0$ ,  $\Delta$  is the Laplace operator,

$$\frac{\partial u_\varepsilon}{\partial N} = \sum_{i,j=1}^n a_{i,j}(x, t) \frac{\partial u_\varepsilon}{\partial x_j} \cos(\bar{n}, \bar{x}_i), \quad \frac{\partial u_\varepsilon}{\partial n} = \sum_{i=1}^n \frac{\partial u_\varepsilon}{\partial x_i} \cos(\bar{n}, \bar{x}_i),$$

$\bar{n}$  is the vector of the external normal to  $\gamma$ , signs "plus" or "minus" mean that the corresponding value is a limit value while a point  $(x, t)$  approaches  $\gamma_T$  inside or outside the domain  $D_T$ ,  $Q$  is a parameter which can equal either 1 or -1.

Let us show that there exists such a parameter  $\varepsilon_0$  that for all  $\varepsilon < \varepsilon_0$  the auxiliary problem (3), (4) is uniquely solvable at any sign of  $Q$ . In this case:

$$\left\| u - \frac{1}{2}(u_\varepsilon^+ + u_\varepsilon^-) \right\|_{W_2^{2,1}(D_T)} \leq C\varepsilon^2, \quad (5)$$

where  $u_\varepsilon^+$  and  $u_\varepsilon^-$  are the solutions of the auxiliary problem (3), (4) at  $Q = 1$  and  $Q = -1$ , respectively.

Let us seek for the solution of problem (3), (4) as power series  $S_1 = \sum_{k=0}^{\infty} \varepsilon^k v_k$  in the domain  $D_T$ ,  $S_2 = \sum_{k=1}^{\infty} \varepsilon^k w_k$  in the domain  $D_T^1$ . Substituting formally  $S_1, S_2$  in (3), (4) and equating the expressions at corresponding power of  $\varepsilon$ , one can obtain a sequence of problems for the determination of  $v_k$  and  $w_k$ :

$$\begin{aligned} Lv_0 &= f, \quad (x, t) \in D_T, & \frac{\partial w_1}{\partial t} - \Delta w_1 &= 0, \quad (x, t) \in D_T^1, \\ v_0(x, 0) &= \varphi(x), \quad x \in D, & w_1(x, 0) &= 0, \quad x \in D_1, \\ v_0 &= 0, \quad (x, t) \in \gamma_T; & \frac{\partial w_1}{\partial n} &= Q \frac{\partial v_0}{\partial N}, \quad (x, t) \in \gamma_T, \\ & & w_1 &= 0, \quad (x, t) \in \Gamma_T; \end{aligned} \quad (6)$$

for  $k \geq 1$ :

$$\begin{aligned} Lv_k &= 0, \quad (x, t) \in D_T, & \frac{\partial w_{k+1}}{\partial t} - \Delta w_{k+1} &= 0, \quad (x, t) \in D_T^1, \\ v_k(x, 0) &= 0, \quad x \in D, & w_{k+1}(x, 0) &= 0, \quad x \in D_1, \\ v_k &= w_k, \quad (x, t) \in \gamma_T; & \frac{\partial w_{k+1}}{\partial n} &= Q \frac{\partial v_k}{\partial N}, \quad (x, t) \in \gamma_T, \\ & & w_{k+1} &= 0, \quad (x, t) \in \Gamma_T. \end{aligned}$$

With such suppositions all the problems in (6) are uniquely solvable and  $v_k \in W_2^{2,1}(D_T)$ ,  $k = 0, 1, \dots$ ,  $w_k \in W_2^{2,1}(D_T^1)$ ,  $k = 1, 2, \dots$ .

**Theorem 1.** *There exists  $\varepsilon_0$  such that for all  $0 < \varepsilon < \varepsilon_0$  the series  $S_1$  and  $S_2$  are absolutely convergent in the spaces  $W_2^{2,1}(D_T)$  and  $W_2^{2,1}(D_T^1)$  correspondingly, and the following equalities take place*

$$u_\varepsilon = S_1, \quad (x, t) \in D_T, \quad u_\varepsilon = S_2, \quad (x, t) \in D_T^1, \quad (7)$$

where  $u_\varepsilon$  is the solution to problem (3), (4).

**Proof.** Using *a priori* estimates for the solutions, traces properties and agreement conditions on  $\gamma_T$ , we have

$$\begin{aligned} \|w_k\|_{W_2^{2,1}(D_T^1)} &\leq C_2 \left\| \frac{\partial w_k}{\partial n} \right\|_{W_2^{\frac{1}{2},1}(\gamma_T)} = C_2 \left\| \frac{\partial v_{k-1}}{\partial N} \right\|_{W_2^{\frac{1}{2},1}(\gamma_T)} \\ &\leq C_2 C_3 \|v_{k-1}\|_{W_2^{2,1}(D_T)}, \end{aligned} \quad (8)$$

where the constants  $C_2$  and  $C_3$  depend only on the regions  $D$ ,  $D_1$  and the coefficients of the operator  $L$  and are  $\varepsilon$ -independent.

The estimate (8) shows that the convergence of the series  $S_1$  in the norm of the space  $W_2^{2,1}(D_T)$  is followed by the convergence of the series  $S_2$  in the norm of the space  $W_2^{2,1}(D_T^1)$ . Let us prove the convergence of the series  $S_1$ .

$$\|v_k\|_{W_2^{2,1}(D_T)} \leq C_4 \|v_k\|_{W_2^{\frac{3}{2},1}(\gamma_T)} = C_4 \|w_k\|_{W_2^{\frac{3}{2},1}(\gamma_T)} \leq C_4 C_5 \|w_k\|_{W_2^{2,1}(D_T^1)}.$$

Taking into account (2) and (8), one can obtain

$$\begin{aligned} \|v_k\|_{W_2^{2,1}(D_T)} &\leq C_6 \|v_{k-1}\|_{W_2^{2,1}(D_T)}, \quad k \geq 1, \\ \|v_0\|_{W_2^{2,1}(D_T)} &\leq C_1 (\|f\|_{L_2(D_T)} + \|\varphi\|_{W_2^{\frac{1}{2}}(D)}), \end{aligned} \quad (9)$$

where  $C_6 = C_2 C_3 C_4 C_5$ .

Thus, if  $\varepsilon < \varepsilon_0 = C_6^{-1}$ , the series  $S_1$  is absolutely convergent in the norm of  $W_2^{2,1}(D_T)$ , and therefore, the series  $S_2$  is also absolutely convergent in the norm of  $W_2^{2,1}(D_T^1)$ . From the completeness of the space  $W_2^{2,1}$  it follows that  $S_1 \in W_2^{2,1}(D_T)$  and  $S_2 \in W_2^{2,1}(D_T^1)$ .

Multiplying equations (6) for  $v_k$  and  $w_k$  by  $\varepsilon^k$  and summing with respect to  $k$ , one can obtain

$$\begin{aligned} LS_1 &= f, \quad (x, t) \in D_T, & \frac{\partial S_2}{\partial t} - \Delta S_2 &= 0, \quad (x, t) \in D_T^1, \\ S_1(x, 0) &= \varphi(x), \quad x \in D, & S_2(x, 0) &= 0, \quad x \in D_1, \\ S_1 &= S_2, \quad (x, t) \in \gamma_T; & \frac{\partial S_2}{\partial n} &= Q\varepsilon \frac{\partial S_1}{\partial N}, \quad (x, t) \in \gamma_T, \\ & & S_2(x, t) &= 0, \quad (x, t) \in \Gamma_T. \end{aligned} \quad (10)$$

Thus,  $u_\varepsilon = S_1$  in  $D_T$ ,  $u_\varepsilon = S_2$  in  $D_T^1$  for all  $0 < \varepsilon < \varepsilon_0$ .  $\square$

It follows from Theorem 1 that the auxiliary problem (3), (4) is uniquely solvable for all  $0 < \varepsilon < \varepsilon_0$  and the following approximation estimate takes place

$$\begin{aligned} \|u - u_\varepsilon^+\|_{W_2^{2,1}(D_T)} &\leq C_7' \varepsilon (\|f\|_{L_2(D_T)} + \|\varphi\|_{\dot{W}_2^1(D)}), \\ \|u - u_\varepsilon^-\|_{W_2^{2,1}(D_T)} &\leq C_7'' \varepsilon (\|f\|_{L_2(D_T)} + \|\varphi\|_{\dot{W}_2^1(D)}), \end{aligned} \quad (11)$$

where  $u_\varepsilon^+ = u_\varepsilon$  at  $Q = 1$ ,  $u_\varepsilon^- = u_\varepsilon$  at  $Q = -1$ . For the proof it is sufficient to compare  $u$  with the expansion  $u_\varepsilon$  in the series  $S_1$  and to take into consideration that the first member of the expansion coincides with  $u$ .

**Theorem 2.** For all  $0 < \varepsilon < \varepsilon_0$  the following estimate takes place

$$\left\| u - \frac{1}{2}(u_\varepsilon^+ + u_\varepsilon^-) \right\|_{W_2^{2,1}(D_T)} \leq C_8 \varepsilon^2 (\|f\|_{L_2(D_T)} + \|\varphi\|_{\dot{W}_2^1(D)}), \quad (12)$$

where  $u$  is the solution of (1),  $u_\varepsilon^+$ ,  $u_\varepsilon^-$  are the solutions to the auxiliary problem (3), (4) at  $Q = 1$  and  $Q = -1$  respectively.

**Proof.** It is obvious from Theorem 1 that the following expansion is true:

$$u_\varepsilon^+ = \sum_{k=0}^{\infty} \varepsilon^k v_k^+, \quad (x, t) \in D_T, \quad u_\varepsilon^+ = \sum_{k=1}^{\infty} \varepsilon^k w_k^+, \quad (x, t) \in D_T^1. \quad (13)$$

where  $v_k^+$ ,  $w_k^+$  are the solutions of (6) at  $Q = 1$ .

Similarly, it is evident from Theorem 1 that the expansion for  $u_\varepsilon^-$  is true

$$u_\varepsilon^- = \sum_{k=0}^{\infty} \varepsilon^k v_k^-, \quad (x, t) \in D_T, \quad u_\varepsilon^- = \sum_{k=1}^{\infty} \varepsilon^k w_k^-, \quad (x, t) \in D_T^1, \quad (14)$$

where  $v_k^-$ ,  $w_k^-$  are the solutions of (6) at  $Q = -1$ .

One can see that  $v_0^+ \equiv v_0^- \equiv u$ , where  $u$  is the solution to the problem (1).

Let us consider  $\tilde{w}_1 = w_1^+ + w_1^-$ . The function  $\tilde{w}_1$  is the solution to the problem

$$\begin{aligned} \frac{\partial \tilde{w}_1}{\partial t} - \Delta \tilde{w}_1 &= 0, \quad (x, t) \in D_T^1, & \frac{\partial \tilde{w}_1}{\partial n} &= 0, \quad (x, t) \in \gamma_T, \\ \tilde{w}_1(x, 0) &= 0, \quad x \in D_1, & \tilde{w}_1 &= 0, \quad (x, t) \in \Gamma_T, \end{aligned}$$

whence it follows that  $\tilde{w}_1 = 0$  or  $w_1^+ = -w_1^-$ .

Then, designate  $\tilde{v}_1 = v_1^+ + v_1^-$ . The function  $\tilde{v}_1$  is the solution to the problem

$$L\tilde{v}_1 = 0, \quad (x, t) \in D_T,$$

$$\tilde{v}_1(x, 0) = 0, \quad x \in D,$$

$$\tilde{v}_1(x, t) = 0, \quad (x, t) \in \gamma_T,$$

whence it appears that  $\tilde{v}_1 = 0$ . It means that  $v_1^+ = -v_1^-$ .

Considering  $\tilde{w}_2 = w_2^+ - w_2^-$ ,  $\tilde{v}_2 = v_2^+ - v_2^-$  one can obtain  $w_2^+ = w_2^-$ ,  $v_2^+ = v_2^-$ . Using this line of reasoning, one can derive the correlation

$$\begin{aligned} v_k^+ &= v_k^-, & \text{when } k \text{ is even,} \\ v_k^+ &= -v_k^-, & \text{when } k \text{ is odd.} \end{aligned} \quad (15)$$

Taking into account (15), one can rewrite the expansions (13) and (14) in the domain  $D_T$

$$\begin{aligned} u_\varepsilon^+ &= u + \varepsilon v_1^+ + \varepsilon^2 v_2^+ + \dots, \\ u_\varepsilon^- &= u - \varepsilon v_1^+ + \varepsilon^2 v_2^+ - \dots \end{aligned} \quad (16)$$

Taking advantage of the representations (16) and the estimate (9), one can find that for all  $0 < \varepsilon < \varepsilon_0$

$$\begin{aligned} \|u - \frac{1}{2}(u_\varepsilon^+ + u_\varepsilon^-)\|_{W_2^{2,1}(D_T)} &\leq \varepsilon^2 \|v_2^+ + \varepsilon^2 v_4^+ + \dots\|_{W_2^{2,1}(D_T)} \\ &\leq C'_8 \varepsilon^2 \|v_0^+\|_{W_2^{2,1}(D_T)} \\ &\leq C_8 \varepsilon^2 (\|f\|_{L_2(D_T)} + \|\varphi\|_{W_2^1(D)}). \quad \square \end{aligned}$$

One can see from the proof of Theorem 2 that main members of errors for  $u_\varepsilon^+$ ,  $u_\varepsilon^-$  in the expansions (16) are of different signs. From the supposition of the function sufficient smoothness it follows that the approximate solutions obtained by means of the fictitious regions method are locally two-sided. It allows to increase the approximation order: the semi-sum of the solutions  $u_\varepsilon^+$ ,  $u_\varepsilon^-$  has a higher order than each of  $u_\varepsilon^+$ ,  $u_\varepsilon^-$ .

## References

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